

1.2. Symmetries and Conservation laws

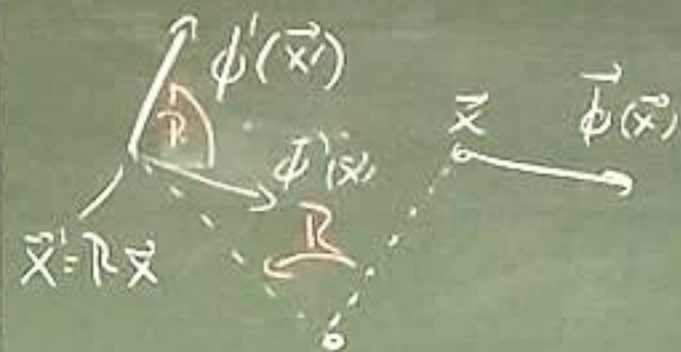
1] $\&$ Transformations of fields decompose into two transformations:

$$x \mapsto x'(x), \quad \phi(x) \mapsto \phi'(x') = \tilde{F}(\phi(x))$$

These are active transformations ∇

Example 1.2 Rotation of vector field $\vec{\phi}$

$$\vec{\phi} = \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}, \quad R \in SO(3)$$



$$\vec{x}' = R\vec{x}$$

$$\vec{\phi}'(\vec{x}') = R\vec{\phi}(\vec{x}) = R\vec{\phi}(R^{-1}\vec{x}')$$

$$\boxed{\vec{\phi}'(\vec{x}) = R\vec{\phi}(R^{-1}\vec{x})}$$

↑
Vector field

$$\phi \mapsto \phi'$$

2] Change of the action under $\phi \rightarrow \phi'$?

$$S' = S[\phi'] = \int d^d x \mathcal{L}(\phi'(x), \partial_\mu \phi'(x))$$

$$= \int d^d x' \mathcal{L}(\phi'(x'), \partial'_\mu \phi'(x'))$$

$$= \int d^d x' \mathcal{L}(\tilde{F}(\phi(x)), \partial'_\mu \tilde{F}(\phi(x)))$$

Substitution $x' = x'(x)$

$$\int d^d x \left| \frac{\partial x'}{\partial x} \right| \mathcal{L}(\tilde{F}(\phi(x)), \frac{\partial x^\alpha}{\partial x'^\mu} \partial_\alpha \tilde{F}(\phi(x)))$$

Example 1.3. Translation

1) $x' = x + a$

$$\phi'(x') = \phi(x) = \phi(x' - a)$$

2) $\mathcal{F} = \mathbb{1}$, $\phi'(x) = \mathcal{F}(\phi(x)) = \phi(x(x'))$

$$\frac{\partial x^\nu}{\partial x'^\mu} = \delta_\mu^\nu$$

3)

$$S[\phi'] = \int d^d x \mathcal{L}(\phi'(x), \partial_\mu \phi'(x)) =$$

$$= \int d^d x \mathcal{L}(\phi(x), \partial_\mu \phi(x)) = S[\phi]$$

→ Action is translation invariant

Example 1.4. Scale transformations

⇒ see script.

Example 1.5. Phase rotation

1) $x' = x$, $\phi'(x') = e^{i\theta} \phi(x)$

2) $\mathcal{F}(\phi) = e^{i\theta} \phi$, $\frac{\partial x^\nu}{\partial x'^\mu} = \delta_\mu^\nu$, $\left| \frac{\partial x'}{\partial x} \right| = 1$

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)(\partial^\mu \phi^*) - \frac{1}{2} m^2 \phi \phi^*$$

Infinitesimal transformations

1) \mathcal{F} Infinitesimal transformations (IT)

$$x'^M = x^M + \omega_a \left[\frac{\delta x^M}{\delta \omega_a} \right]$$

$$\phi'(x') = \phi(x) + \omega_a \frac{\delta \mathcal{F}}{\delta \omega_a}(x) = \mathcal{F}(\phi(x))$$

For now: $\omega_a = \omega_a(x)$

2) Generator of IT:

$$\delta_\omega \phi(x) = \phi'(x) - \phi(x) \equiv -i \omega_a \underline{G_a \phi(x)}$$

with

$$\phi'(x) = \phi(x) - \omega_a \frac{\delta x^a}{\delta \omega_c} \partial_\mu \phi(x)$$

$$+ \omega_a \frac{\delta \mathcal{F}}{\delta \omega_a}(x) + O(\omega^2)$$

$$\Rightarrow i G_a \phi = \frac{\delta x^a}{\delta \omega_a} \partial_\mu \phi - \frac{\delta \mathcal{F}}{\delta \omega_a}$$

Example 16. Translation

$$1) x'^\mu = x^\mu + \omega^\mu = x^\mu + \omega^\nu \left[\frac{\delta x^\mu}{\delta \omega^\nu} \right]$$

$$2) \frac{\delta \mathcal{F}}{\delta \omega^\nu} = 0 \quad (\text{Scalar field}) \quad \delta_\omega^\mu = 0$$

$$3) i G_\mu \phi = \delta_\mu^\nu \partial_\nu \phi - 0$$

$$\rightarrow G_\mu = -i \partial_\mu = P_\mu \quad \left[\begin{array}{l} -i \partial_t \psi = H \psi \\ \underline{\underline{\hspace{1cm}}} \end{array} \right]$$

$$S[\phi'] = S[\phi] + \int d^4x \partial_\mu K^\mu$$

Ex. 17

Ex. 18. Spatial Rotations

$$G_{\mu\nu} = i(x_\nu \partial_\mu - x_\mu \partial_\nu) + \underline{S_{\mu\nu}}$$

Noether's theorem

1) Transformations for which

$$S[\phi] = S[\phi'] \Leftrightarrow \begin{array}{l} \text{(strict)} \\ \text{symmetry of} \\ \text{the action} \end{array}$$

for ω_a independent of x (rigid transform.)

2] Assume in (*) $\omega_a = \omega_a(x)$ is not rigid

$$\uparrow O(\omega_a) = O(\partial_\mu \omega_a)$$

$$3] \frac{\partial x'^\mu}{x^\mu} = \delta_\mu^0 + \partial_\mu \left(\omega_a \frac{\delta x^\mu}{\delta \omega_a} \right)$$

$$\cdot \left| \frac{\partial x'}{\partial x} \right| = 1 + \partial_\mu \left(\omega_a \frac{\delta x^\mu}{\delta \omega_a} \right)$$

$$\left[\det(1+A) = 1 + \text{Tr}[A] + O(A^2) \right]$$

4] Inverse Jacobian matrix.

$$\frac{\partial x^0}{\partial x'^\mu} = \delta_\mu^0 - \left(\partial_\mu \left(\omega_a \frac{\delta x^0}{\delta \omega_a} \right) \right)$$

$$5] S' = S[\phi'] = \int d^4x \left[1 + \partial_\mu \left(\omega_a \frac{\delta x^\mu}{\delta \omega_a} \right) \right]$$

$$\times \left(\phi + \omega_a \frac{\delta \mathcal{F}}{\delta \omega_a} \left[\delta_\mu^0 - \partial_\mu \left(\omega_a \frac{\delta x^0}{\delta \omega_a} \right) \right] \right)$$

$$\times \left[\partial_0 \phi + \partial_0 \left(\omega_a \frac{\delta \mathcal{F}}{\delta \omega_a} \right) \right]$$

6] Expand in 1st order of $\omega_a, \partial_\mu \omega_a$

$$7] \delta S = S' - S = \int d^4x \left[\underbrace{K^a}_{\text{Symmetry}} \omega_a - \int_a^M \partial_\mu \omega_a \right]$$

Symmetry
 $(\Rightarrow \omega_a \text{ rigid})$

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)(\partial^\mu \phi^*) - \frac{1}{2} m^2 \phi^*$$

Infinitesimal transformations

1] \mathcal{L} Infinitesimal transformations (IT)

$$\begin{cases} x'^\mu = x^\mu + \underbrace{[\omega_a]}_{\text{matrix}} \left[\frac{\delta x^\mu}{\delta \omega_a}(x) \right] \\ \phi'(x') = \phi(x) + \omega_a \frac{\delta \mathcal{F}}{\delta \omega_a}(x) = \mathcal{F}(\phi(x)) \end{cases}$$

For non. $\omega_a = \omega_a(x)$

2) For non-rigid transformation

$$\delta S = - \int d^d x j_a^\mu \partial_\mu \omega_a$$

with the current:

$$j_a^\mu = \left\{ \frac{\partial \mathcal{L}}{\partial(\partial_\mu \psi)} \partial_\mu \psi - \delta_a^\mu \mathcal{L} \right\} \frac{\delta x^0}{\delta \omega_a} - \frac{\partial \mathcal{L}}{\partial(\partial_\mu \psi)} \frac{\delta \mathcal{F}}{\delta \omega_a}$$

ϕ that obey the EOM

$$\Rightarrow \delta S = 0 \text{ for arbitrary variations } \phi' = \phi + \delta \phi$$

$$\Rightarrow \text{In particular for non-rigid transformations}$$

$$\partial_\mu j_a^\mu = 0 \quad \forall x, a$$

Noether's (first) theorem

9) Integration by parts:

$$\delta S = \int d^d x \omega_a(x) \partial_\mu j_a^\mu = 0$$

11) Conserved charge:

$$Q_a = \int_{\text{Space}} d^{d-1} x j_a^0$$

Note 1.3:

$$P_{\mu\nu} = \frac{1}{2} g_{\mu\nu} P = \frac{8\pi G}{c} T_{\mu\nu}$$

$$T^{\mu\nu} \neq T^{\nu\mu}$$

$$\text{But } T^{\mu\nu} = T^{\nu\mu} + \partial_\rho h^{\mu\nu\rho}$$

Homework:

$$\frac{dQ_a}{dt} = \dots = - \int_{\text{Surface}} \text{div } j_a^\mu = 0$$

Note 1.1 j_a^μ canonical current

$$\tilde{j}_a^\mu = j_a^\mu + \partial_\nu B_a^{\mu\nu}, \quad B_a^{\mu\nu} = -B_a^{\nu\mu}$$

arbitrary

$$\Rightarrow \partial_\mu \tilde{j}_a^\mu = 0$$

Note 1.2:

$$\text{Symmetric } \mathcal{L} \Rightarrow \text{Symmetric } S \Rightarrow \text{Symmetric EOM}$$

\rightarrow conserved charges

$$\mathcal{L} = \frac{1}{2} (\partial\phi)^2 - \frac{1}{2} \phi^2 \Rightarrow (\partial^2 + m^2)\phi = 0$$

$$\phi' = \lambda\phi \quad SO(1,3) \times \mathbb{R}^4$$

Application: Energy-Momentum Tensor (EMT)

1) Lorentz spacetime. $x'^M = x^M + \xi^M$

$$\rightarrow \frac{\delta x^M}{\delta \xi^0} = \delta^M_0$$

$$\frac{\delta \tilde{\mathcal{L}}}{\delta \xi^0} = 0$$

2) Translation invariant action S

3) Conserved currents

$$T^M_{\nu} = \left\{ \frac{\partial \mathcal{L}}{\partial (\partial_\nu \phi)} \partial_\nu \phi - \delta^M_\nu \mathcal{L} \right.$$

$$T^{\mu 0} = g^{\nu\rho} T^{\mu}_{\rho} = \frac{\partial \mathcal{L}}{\partial (\partial_\nu \phi)} \partial^\nu \phi - g^{\mu\nu} \mathcal{L}$$

Energy-momentum tensor.

with $\partial_\mu T^{\mu 0} = 0$

$$P^0 = \int d^3x T^{00}$$

Conserved charges

4) Energy ($\nu=0$).

$$P^0 = \int d^3x T^{00} = \int d^3x \left\{ \underbrace{\frac{\partial \mathcal{L}}{\partial \dot{\phi}} \dot{\phi}}_{\pi} - \mathcal{L} \right\}$$

= \mathcal{H} Hamiltonian \mathcal{H}

5) $\nu=i \Rightarrow$ Kinetic momentum

$$P^i = \int d^3x T^{0i} = \int d^3x \frac{\partial \mathcal{L}}{\partial \dot{\phi}} (-\partial_i \phi)$$

$$= - \int d^3x \pi (\partial_i \phi)$$