

## 2. The Klein-Gordon Field

### 2.1. Canonical Quantization

#### 1. Theory.

$$\phi = \phi^*$$

i) Real field  $\phi(x)$

ii) Lagrangian:  $\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{1}{2} m^2 \phi^2$

iii) EOM:  $(\partial^2 + m^2) \phi = 0$

iv) Hamiltonian:  $\mathcal{H} = \frac{1}{2} \pi^2 + \frac{1}{2} (\nabla \phi)^2 + \frac{1}{2} m^2 \phi^2$

### 2. Canonical quantization.

$$[\phi(\vec{x}), \phi(\vec{y})] = 0$$

$$[\pi(\vec{x}), \pi(\vec{y})] = 0$$

$$[\phi(\vec{x}), \pi(\vec{y})] = i \delta^{(3)}(\vec{x} - \vec{y})$$

(3)

$$\phi^\dagger = \phi, \quad \pi^\dagger = \pi$$

### 3. Goals:

- Representation of field operators
- Spectrum of Hamiltonian
- Time evolution

### 4. Motivation.

i) FT UG equation in space.

$$\phi(\vec{x}, t) = \int \frac{d^3 p}{(2\pi)^3} e^{i\vec{p}\vec{x}} \tilde{\phi}(\vec{p}, t)$$

$$\text{UG} \rightarrow (\partial_t^2 + \underbrace{|\vec{p}|^2 + m^2}_{\omega_p^2}) \tilde{\phi}(\vec{p}, t) = 0$$

→ Decoupled Harmonic oscillators.

and constraint  $\phi^*(\vec{p}, t) = \phi(-\vec{p}, t)$

$$\text{ii) } \mathcal{H}_{\text{SHO}} = \frac{1}{2} \tilde{\pi}^2 + \frac{1}{2} \omega_p^2 \tilde{\phi}^2$$

$$\tilde{\phi} = \frac{1}{\sqrt{2\omega}} (a + a^\dagger)$$

$$\tilde{\pi} = -i\sqrt{\frac{\omega}{2}} (a - a^\dagger)$$

$$\text{with } [a, a^\dagger] = 1$$

$$\Rightarrow \mathcal{H}_{\text{SHO}} = \omega (a^\dagger a + \frac{1}{2})$$

### 5. Field operators.

$$\phi(\vec{x}) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2\omega_{\vec{p}}}} (a_{\vec{p}} + a_{-\vec{p}}^\dagger) e^{i\vec{p}\cdot\vec{x}}$$

$$= \int \frac{d^3p}{(2\pi)^3} (a_{\vec{p}} e^{i\vec{p}\cdot\vec{x}} + a_{-\vec{p}}^\dagger e^{-i\vec{p}\cdot\vec{x}}) \quad (1)$$

$$\pi(\vec{x}) = \int \frac{d^3p}{(2\pi)^3} -i\sqrt{\frac{\omega_{\vec{p}}}{2}} (a_{\vec{p}} - a_{-\vec{p}}^\dagger) e^{i\vec{p}\cdot\vec{x}} \quad (2)$$

$$(\phi^\dagger = \phi, \pi^\dagger = \pi)$$

with momentum modes

$$[a_{\vec{p}}, a_{\vec{q}}^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{q}) \quad (2)$$

(check: (1) + (2)  $\Rightarrow$  (3))  $\omega_{\vec{p}} = \sqrt{\vec{p}^2 + m^2} = E_{\vec{p}}$

### 6. Hamiltonian.

$$H = \int \frac{d^3p}{(2\pi)^3} \omega_{\vec{p}} (a_{\vec{p}}^\dagger a_{\vec{p}} + \frac{1}{2} [a_{\vec{p}} a_{\vec{p}}^\dagger])$$

$$\frac{1}{2} [a_{\vec{p}} a_{\vec{p}}^\dagger] \rightarrow \frac{1}{2} \delta(0) = \infty$$

( $\rightarrow$  Normal ordering)

### 7. Eigenstates + Spectrum.

i)  $\rightarrow [H, a_{\vec{p}}^\dagger] = \omega_{\vec{p}} a_{\vec{p}}^\dagger$

ii) Vacuum  $|0\rangle \rightarrow$  Eigenstates

$$a_{\vec{p}} |0\rangle = 0 \quad \forall \vec{p} \quad a_{\vec{p}}^\dagger a_{\vec{q}}^\dagger \dots |0\rangle$$

iii) Energy:  $E_{\vec{p}} = \omega_{\vec{p}} = \sqrt{\vec{p}^2 + m^2}$

iv) Kinetic momentum

$$P^i = \int d^3x \pi(\vec{x}) (-\partial_i \phi(\vec{x}))$$

$$= \int \frac{d^3p}{(2\pi)^3} p_i a_{\vec{p}}^\dagger a_{\vec{p}}$$

v) Statistics.  $a_{\vec{p}}^\dagger a_{\vec{q}}^\dagger |0\rangle = a_{\vec{q}}^\dagger a_{\vec{p}}^\dagger |0\rangle$

$\rightarrow$  Excitations  $a_{\vec{p}}^\dagger$  commute and carry additive energy + momentum

$\rightarrow$  Bosonic particles

8. Normalization

$$i) \Lambda = R L_3(p) R \in SO^+(1,3)$$

$$P' = (\underline{E}_{P'}, \vec{P}') = \Lambda R (E_P, \vec{P})$$

ii) Jacobian in space

$$\det \left( \frac{\partial \vec{P}'}{\partial \vec{P}} \right) = \frac{dP'_3}{dP_3} = \frac{E_{P'}}{E_P}$$

$$\rightarrow \delta^{(3)}(\vec{P}' - \vec{q}') = \frac{E_{P'}}{E_P} \delta^{(3)}(\vec{P} - \vec{q})$$

not LI

$\rightarrow E_P \delta(\vec{P} - \vec{q})$  is LI

iii) Single particle eigenstates

$$|\vec{P}\rangle = \sqrt{E_P} a_P^\dagger |0\rangle$$

$$\Rightarrow \langle \vec{P} | \vec{q} \rangle = (2\pi)^3 2 E_P \delta(\vec{P} - \vec{q})$$

9. Lorentz transformations  $\Lambda \in SO^+(1,3)$

$$U(\Lambda) |\vec{P}\rangle := |\Lambda \vec{P}\rangle$$

$$\Leftrightarrow U(\Lambda) a_{\vec{P}}^\dagger U^\dagger(\Lambda) = \sqrt{\frac{E_{\Lambda \vec{P}}}{E_{\vec{P}}}} a_{\Lambda \vec{P}}^\dagger$$

10. Interpretation of  $\phi(x)$

$$\phi(\vec{x}) |0\rangle = \int \frac{d^3P}{(2\pi)^3} \frac{1}{2E_P} e^{-i\vec{P}\vec{x}} |\vec{P}\rangle$$

$|\vec{P}| \ll m \Rightarrow E_P \approx \text{const.}$

$\rightarrow \phi(\vec{x})$  creates particle at  $\vec{x}$

$$(\langle 0 | \phi(x) | \vec{P} \rangle = e^{i\vec{P}\vec{x}})$$

Note 2.1:

• Projector on SP sector

$$1_1 = \int \frac{d^3P}{(2\pi)^3} |\vec{P}\rangle \frac{1}{2E_P} \langle \vec{P}|$$

•  $\oint \frac{d^3P}{(2\pi)^3} \frac{1}{2E_P} f(P)$  (LI)  $f(P) = f(P')$

$$\Rightarrow \text{LI}$$

# 2.2. The Klein-Gordon Field in Space-Time

1. Heisenberg operators:

$$\phi(\vec{x}, t) = \phi(\vec{x}) = e^{iHt} \phi(\vec{x}) e^{-iHt}$$

2. Heisenberg equation:

$$i\partial_t \phi = [\phi, H]$$

•  $i\partial_t \phi(x) = \left[ \phi(x), \int d^3y \left\{ \frac{1}{2} \pi^2 + \frac{1}{2} (\nabla \phi)^2 + \frac{1}{2} m^2 \phi^2 \right\} \right]$

$\stackrel{0}{=} i\pi(x)$        $\pi(x) = \partial_t \phi(x)$

•  $i\partial_t \pi(x) \stackrel{0}{=} -i(-\nabla^2 + m^2)\phi(x)$  (\*\*)

(\*)  $\Rightarrow (\partial_t^2 - \nabla^2 + m^2)\phi(x) = 0$  *Klein Gordon equation*

3. Time evolution of modes

$$e^{iHt} a_{\vec{p}} e^{-iHt} = a_{\vec{p}} e^{\pm iE_{\vec{p}}t}$$

4. Field operators:

$$\phi(x) = \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_{\vec{p}}}} \left( a_{\vec{p}} e^{-ipx} + a_{\vec{p}}^\dagger e^{ipx} \right)$$

$\pi(x) = \partial_t \phi(x)$

$x-y \geq 0$   $\langle 0 | \phi(x) \phi(y) | 0 \rangle$

$p^0 = E_{\vec{p}}$

Note 2.2

1.  $\phi(x, t) = e^{iHt} \phi(\vec{x}, 0) e^{-iHt}$   
(Translation in time)

•  $b(\vec{x}) = e^{-i\vec{P}\vec{x}} \phi_s(0) e^{i\vec{P}\vec{x}}$

$\Rightarrow \phi(x) = e^{iP^M x_M} \phi(0) e^{-iP^M x_M}$

Note 2.3

•  $e^{-iPx}$   $\rightarrow$  pos. freq. solutions of KG equation

•  $e^{+iPx}$   $\rightarrow$  annihilation of  $a_{\vec{p}}$  creation  $a_{\vec{p}}^\dagger$